

**The influence of a new band gap narrowing model on measurements of the intrinsic carrier density in crystalline silicon**

Pietro P. Altermatt,<sup>a)</sup> Andreas Schenk,<sup>b)</sup> Gernot Heiser,<sup>c),a)</sup> Martin A. Green<sup>d)</sup>

- a) University of NSW, Photovoltaic Special Research Centre, Sydney 2052, Australia, pietro@cse.unsw.edu.au
- b) Integrated Systems Laboratory, ETH Zurich, Gloriastr. 35, 8092 Zurich, Switzerland, schenk@iis.ee.ethz.ch
- c) University of NSW, School of Computer Science&Engineering, Sydney 2052, Australia, G. Heiser@unsw.edu.au
- d) University of NSW, Photovoltaic Special Research Centre, Sydney 2052, Australia, m.green@unsw.edu.au

**ABSTRACT**

The commonly used value of the intrinsic carrier density of crystalline silicon,  $n_i = 1.00(\pm 0.015) \times 10^{10} \text{ cm}^{-3}$  (at 300 K), was determined by Sproul&Green in 1991 (the value in brackets is the one-standard-deviation uncertainty). This value is the average taken from various devices with differing doping densities in the range of  $1 \times 10^{14}$  to  $1 \times 10^{16} \text{ cm}^{-3}$ . Sproul&Green obtained values of  $n_i$  that increase slightly with increasing doping density. We demonstrate that their experiment was influenced by band gap narrowing, even though the doping density of their samples was low ( $< 1 \times 10^{16} \text{ cm}^{-3}$ ). Our re-interpretation of their measurements leads to results consistent with  $n_i = 9.7(\pm 0.1) \times 10^9 \text{ cm}^{-3}$ , recently determined by Misiakos&Tsamakis.

**1. The  $n_i$  measurement of Sproul&Green**

In 1991, Sproul&Green [1,2] determined the intrinsic carrier density of crystalline silicon,  $n_i$ , to be  $1.00(\pm 0.015) \times 10^{10} \text{ cm}^{-3}$  at 300 K. Nowadays, this is the most widely accepted value of  $n_i$  in the silicon community. In their experiment, Sproul&Green designed solar cells such that their current is limited by the electron flow through the p-type base. Thus, the current-voltage (I-V) curves of these cells depend mainly on  $n_i$ , the mobility of minority electrons in the base,  $\mu$ , the boron doping density in the base,  $N_A$ , and the base thickness,  $W$ . Sproul&Green obtained  $n_i$  by reproducing their measured I-V curves with an analytical model. Only  $n_i$  had to be adjusted as a free model parameter because  $\mu$ ,  $N_A$ , and  $W$  were known from separate measurements. The symbols in Figure 1 show the measured dark I-V curves [3], and the symbols in Figure 2 depict  $n_i$  extracted from these cells.

**2. Absolute and effective intrinsic carrier density**

Figure 2 suggests that the extracted  $n_i$  increases with increasing doping density. In the original work [2], it is mentioned that band gap narrowing (BGN) may be the cause for this. However, as there were no measurements available

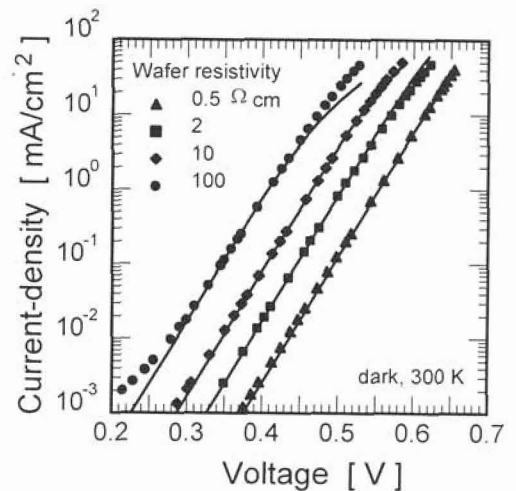


Figure 1: Dark I-V curves, measured by Sproul&Green (symbols) [2], and simulated using Dessis [6] with the BGN model of Schenk [5] (lines).

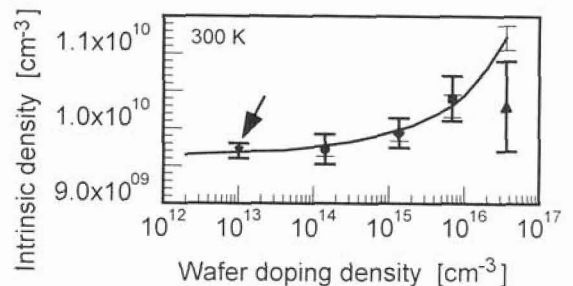


Figure 2: The (effective) intrinsic carrier density of silicon, determined by Sproul&Green with their analytical model [1] (symbols and bold error bars), the value given by Misiakos&Tsamakis [4] (arrow), and simulations with Dessis and the BGN model of Schenk [5] (line and thin error bars).

indicating BGN at  $N_A$  below  $1 \times 10^{17} \text{ cm}^{-3}$ , the average of these  $n_i$  values was taken, ie.  $n_i = 1.00(\pm 0.015) \times 10^{10} \text{ cm}^{-3}$ . In this paper we introduce a model that predicts a considerable amount of BGN at  $N_A > 1 \times 10^{13} \text{ cm}^{-3}$ . This implies that the

values of Sproul&Green in Figure 2 do not represent  $n_i$ , but the effective intrinsic carrier density,  $n_{i,eff}$ , according to the formula

$$n_{i,eff} = n_i \exp[\Delta E_g / 2kT]. \quad (1)$$

With this interpretation, a discrepancy with the measurement of Misiakos&Tsamakis [4] is resolved, who measured  $n_i$  at carrier densities around  $1 \times 10^{13} \text{ cm}^{-3}$  and obtained a lower value of  $9.7(\pm 0.1) \times 10^9 \text{ cm}^{-3}$ . Their value is included in Figure 2 (arrow) and is consistent with the Sproul&Green values at low  $N_A$ . In this paper, we use the BGN model of Schenk [5], who recently developed a band gap narrowing model that is entirely derived from quantum mechanical principles and thus, in contrast to empirical BGN models, is able to calculate BGN at  $N_A$  below  $1 \times 10^{17} \text{ cm}^{-3}$ . We implemented this model into the device simulator Dessis [6], which numerically solves the fully coupled set of semiconductor differential equations. We reproduce the dark I-V curves of Figure 1 with a simulation approach that includes the entire volume of the  $2 \times 2 \text{ cm}^2$  cells, including their perimeter region [7].

### 3. Simulation results

The lines in Figures 1 and 2 show the simulation results. The simulations deviate from the I-V measurements at high and low current-densities of Figure 1 for reasons that are well understood and explained in Ref. 1. The simulated error bars in Figure 2 (thin lines) are obtained by varying  $\mu$  within the uncertainty limits of Sproul&Green's photoconductance decay measurements of  $\mu$  [2] (symbols in Figure 3). The simulated error bars in Figure 2 (thin lines) are smaller than in the original work of Sproul&Green (bold lines in Figure 2) [1,2]. This is so because  $\mu$  is the only parameter that is varied in our simulations, while Sproul&Green considered the uncertainties of all the other parameters as well. As our simulations imitate the measurements, they cannot reduce the uncertainties. Hence, the one-standard-deviation uncertainty of  $n_i$  used in our simulations is at least 3%, as given by Sproul&Green.

As our simulated  $n_{i,eff}$  values reproduce the ' $n_i$  values' of Sproul&Green very well (by choosing a single value of  $n_i = 9.65 \times 10^9 \text{ cm}^{-3}$ ), we conclude that Sproul&Green measured  $n_{i,eff}$  instead of  $n_i$ . This implies that the data of Sproul&Green and of Misiakos&Tsamakis are consistent.

### 4. Recommendation for $n_i$ value

In the interpretation given here,  $n_i$  is the asymptotic value at the left part of Figure 2, ie.  $n_i = 9.65(\pm 0.15) \times 10^9 \text{ cm}^{-3}$ . In most cases, it does not make a substantial difference whether  $n_i = 9.65(\pm 0.15) \times 10^9 \text{ cm}^{-3}$  or  $1.00(\pm 0.015) \times 10^{10} \text{ cm}^{-3}$  is used. These values differ only by 4%. If calculations are made with a BGN model that yields no gap shrinkage below

a dopant density of  $1 \times 10^{17} \text{ cm}^{-3}$ , and the behaviour of the device is ruled by a region with a resistivity of 1-10  $\Omega\text{cm}$ , one may use  $n_i = 1.00(\pm 0.015) \times 10^{10} \text{ cm}^{-3}$ , as this value is the actual  $n_{i,eff}$  of this region. This is the case for many types of solar cells, and explains why  $n_i = 1.00 \times 10^{10} \text{ cm}^{-3}$  has been so successfully applied even in highly accurate simulations [7]. If the BGN model of Schenk is applied, we recommend using  $n_i = 9.65(\pm 0.15) \times 10^9 \text{ cm}^{-3}$ . In cases where the behaviour of the device is ruled by highly doped regions, care must be taken that the applied BGN model was derived with a similar  $n_i$  value as is used in the simulations. Otherwise, the BGN model has to be adjusted to the new  $n_i$  value. For more details, please refer to the paper of Schumacher *et al.* of this conference.

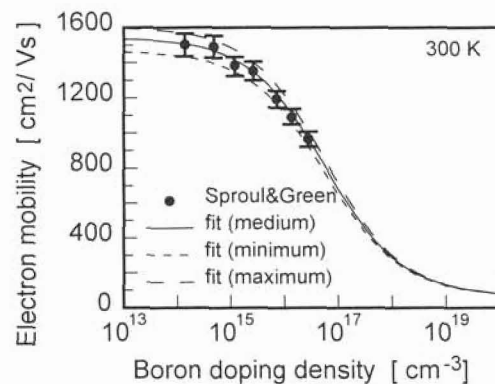


Figure 3: The minority electron mobility, measured by Sproul&Green [2] (symbols), and three parameterisations used in our Dessis model (lines).

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